

An Introduction to Closure Phases

John Monnier

**Interferometry Summer School
August 1999**

Outline of Tutorial

A. Introduction

**B. Closure Phases and
Amplitudes**

C. What to do with Φ_{CP} ?

D. Mapping with Φ_{CP}

i.e. self-calibration

E. New Developments

DETECTION OF ELECTRIC FIELD AT
EACH TELESCOPE INTRODUCES A
COMPLEX GAIN \tilde{G} :

$$\begin{aligned}\tilde{E}_i^{\text{MEASURED}} &= \tilde{G}_i \cdot \tilde{E}_i^{\text{TRUE}} \\ &= |G_i| e^{i\phi_i} \tilde{E}_i^{\text{TRUE}}\end{aligned}$$

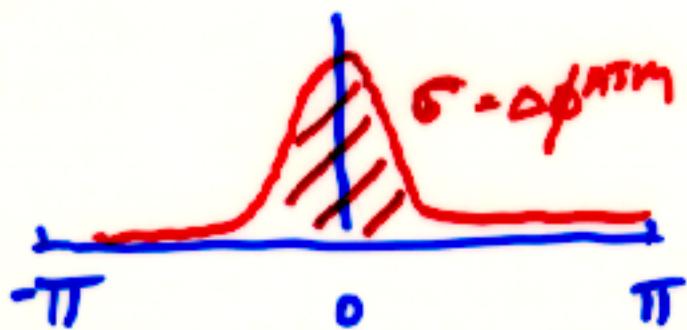
THIS AFFECTS MEASUREMENT OF
THE COMPLEX VISIBILITY $\tilde{\gamma}$:

$$\begin{aligned}\tilde{\gamma}_{ij} &\propto \tilde{E}_i \cdot \tilde{E}_j^* \quad \text{Hence} \\ \tilde{\gamma}_{ij}^{\text{measured}} &= \tilde{G}_i \tilde{G}_j^* \tilde{\gamma}_{ij}^{\text{TRUE}} \\ &= |G_i| |G_j| e^{i(\phi_i - \phi_j)} \tilde{\gamma}_{ij}^{\text{TRUE}}\end{aligned}$$

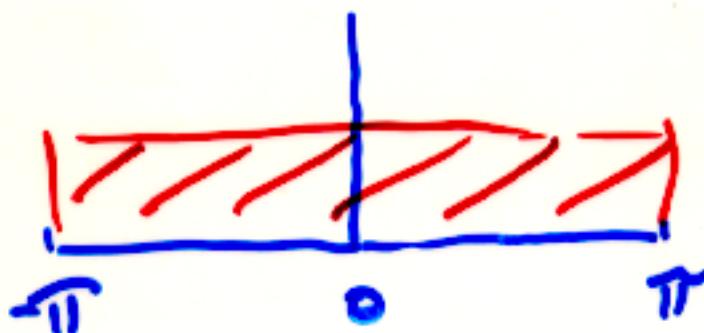
WHY NOT JUST AVERAGE THE SEPARATE PHASES?

$$\phi^{\text{measured}} = \phi^{\text{true}} + \Delta\phi^{\text{ATM}}$$

CASE 1: $\Delta\phi^{\text{ATM}} < \pi$



CASE 2: $\Delta\phi^{\text{ATM}} > \pi$ Then



NO PHASE INFORMATION AT ALL

**Fourier Phases are crucial for
image reconstructions of all but the
simplest sources. What to do?**

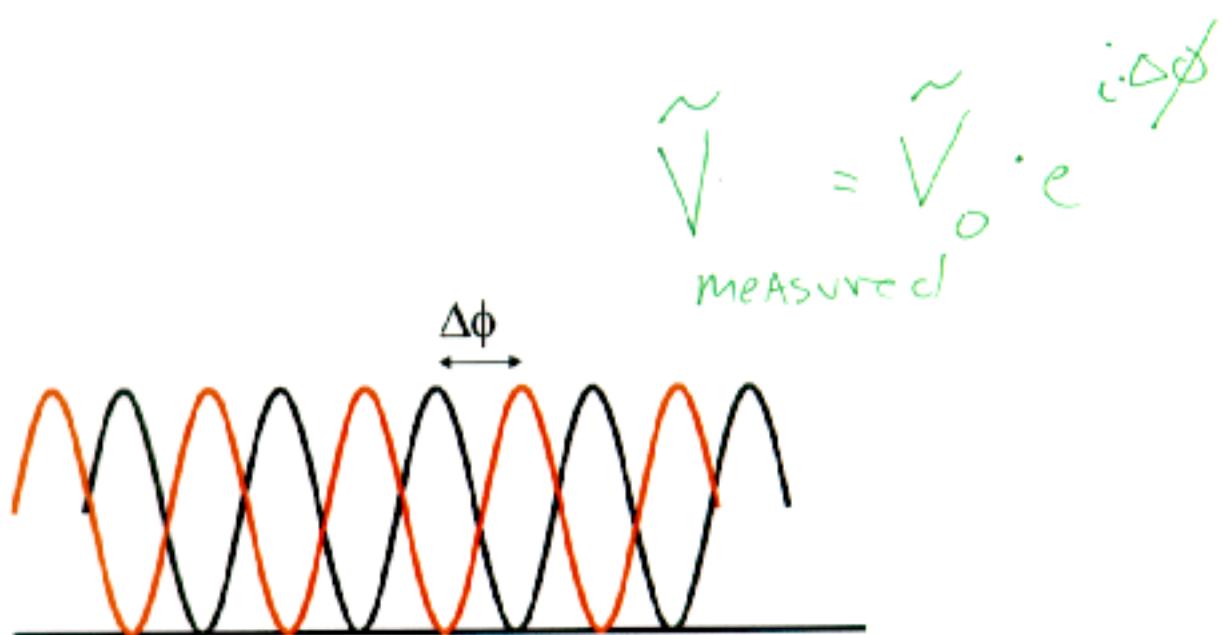
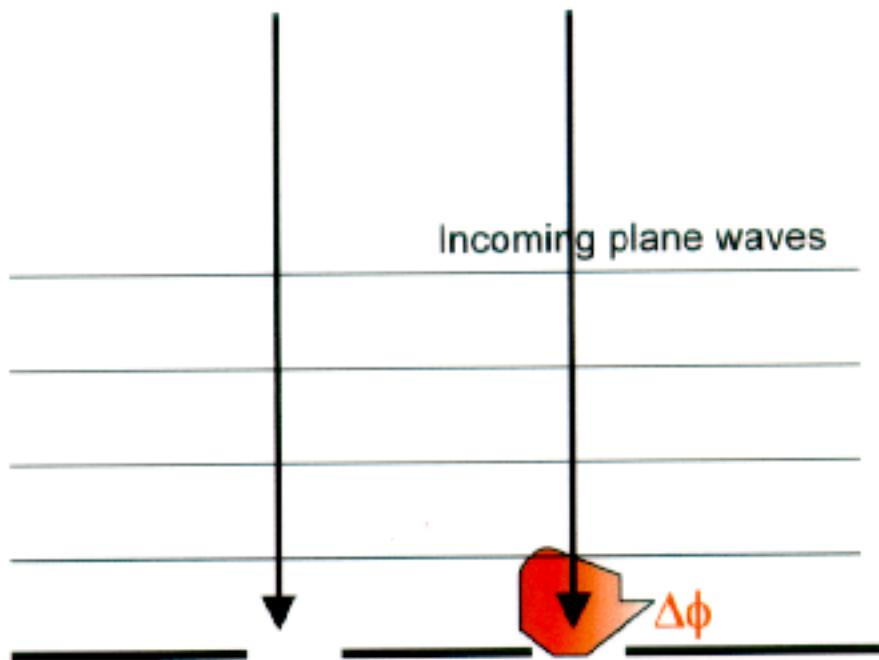
Phase Referencing

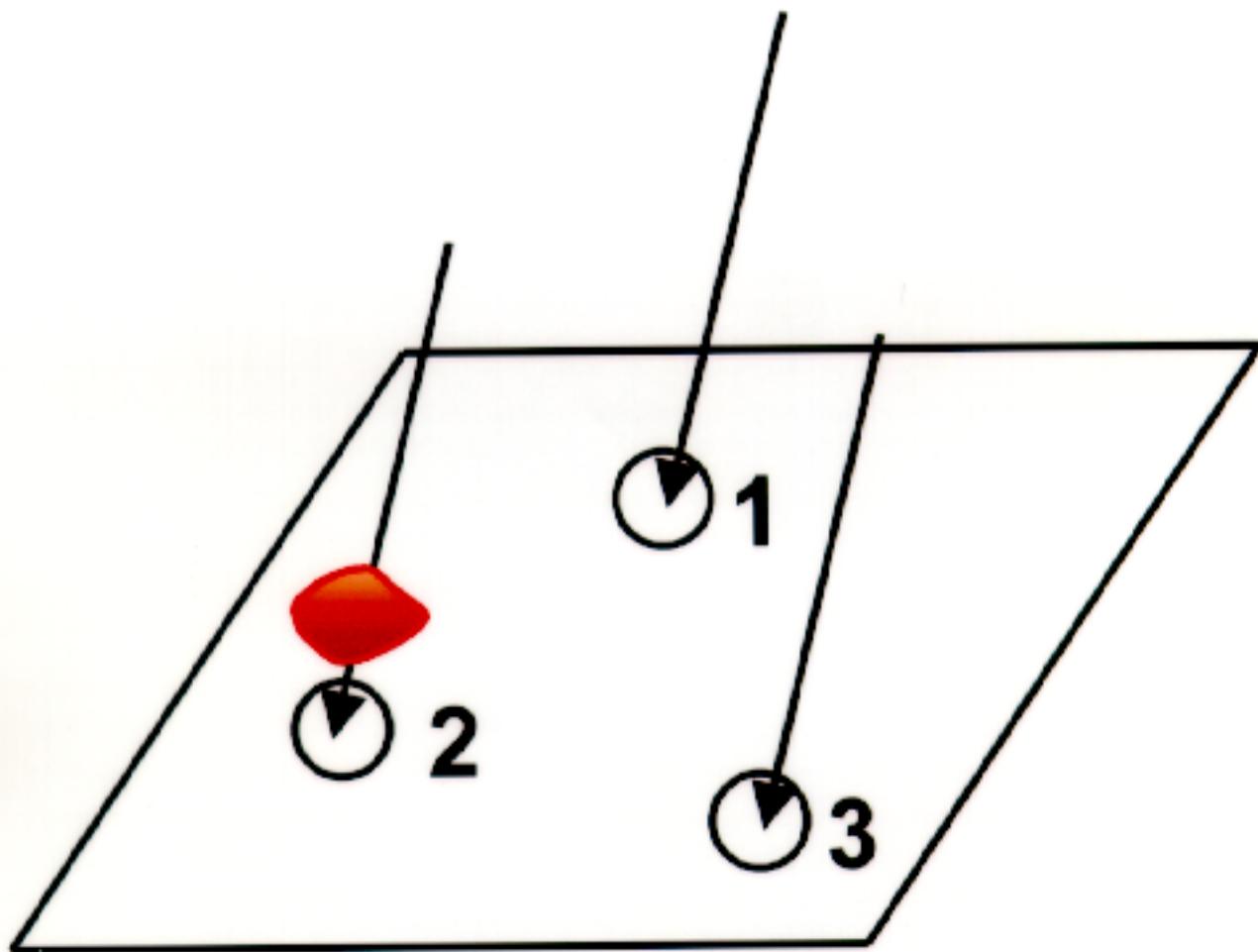
- nearby sources
- measuring $\Delta\phi_{\text{atm}}$
- spectral lines

or

Closure Phases

• Point source
at infinity





Observed Intrinsic Atmosphere

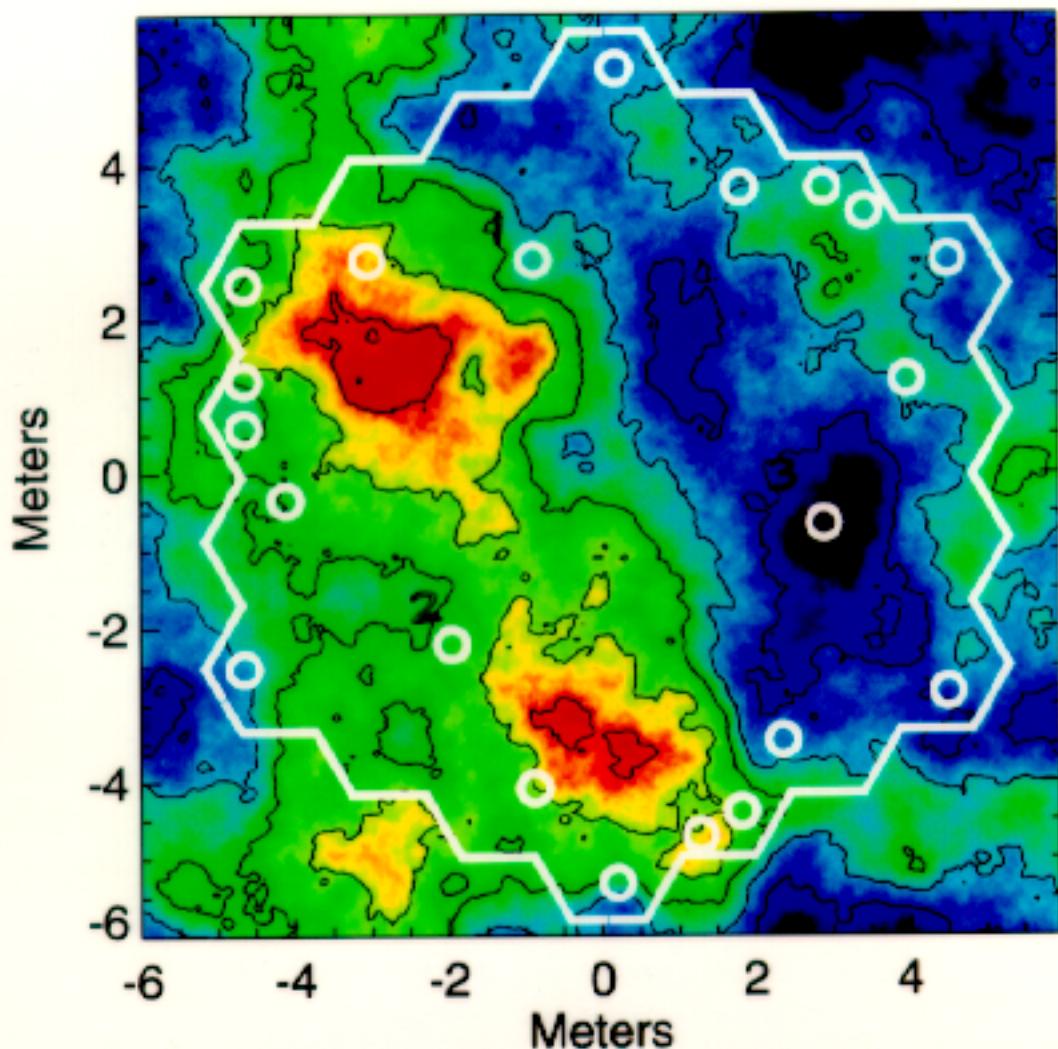
$$\Phi(1-2) = \Phi_o(1-2) + [\phi(2)-\phi(1)]$$

$$\Phi(2-3) = \Phi_o(2-3) + [\phi(3)-\phi(2)]$$

$$\Phi(3-1) = \Phi_o(3-1) + [\phi(1)-\phi(3)]$$

Closure	$= \Phi_o(1-2) + \Phi_o(2-3)$
Phase	$+ \Phi_o(3-1)$
(1-2-3)	

Phasescreen



Observed Intrinsic Atmosphere

$$\Phi(1-2) = \Phi_o(1-2) + [\phi(2)-\phi(1)]$$

$$\Phi(2-3) = \Phi_o(2-3) + [\phi(3)-\phi(2)]$$

$$\Phi(3-1) = \Phi_o(3-1) + [\phi(1)-\phi(3)]$$

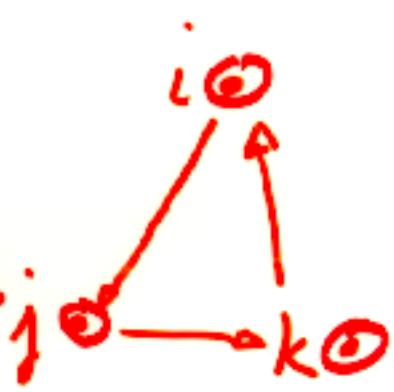
Closure Phase (1-2-3) = $\Phi_o(1-2) + \Phi_o(2-3) + \Phi_o(3-1)$

THE BISPECTRUM

TRIPLE PRODUCTS OF
COMPLEX VISIBILITIES

AROUND CLOSED TRIANGLES

$$\tilde{B}_{ijk} = \hat{\gamma}_{ij} \hat{\gamma}_{jk} \hat{\gamma}_{ki}$$



Hence, AMPLITUDE OF BISPECTRUM

$$|\tilde{B}_{ijk}| = |\tilde{\gamma}_{ij}| |\tilde{\gamma}_{jk}| |\tilde{\gamma}_{ki}|$$

AND PHASE OF BISPECTRUM

$$= \Phi_{ij} + \Phi_{jk} + \Phi_{ki}$$

$$= \Phi_{ijk}^{CP} \rightarrow \text{THE CLOSURE PHASE}$$

General Statements

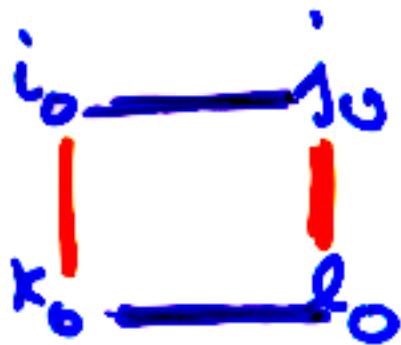
**The Bispectrum is always REAL
for sources with *point-symmetry*
(i.e. the closure phases are 0° or 180°)**

**Closure Phases are not sensitive to an
overall translation of image (a translation
is indistinguishable from atmospheric
phase delays for any given closing
triangle)**

CLOSURE AMPLITUDES

RECALL:

$$\tilde{\gamma}_{ij} = |G_i| |G_j| e^{j(\phi_i - \phi_j)} \\ \times \tilde{\gamma}_{ij}^{\text{TRUE}}$$



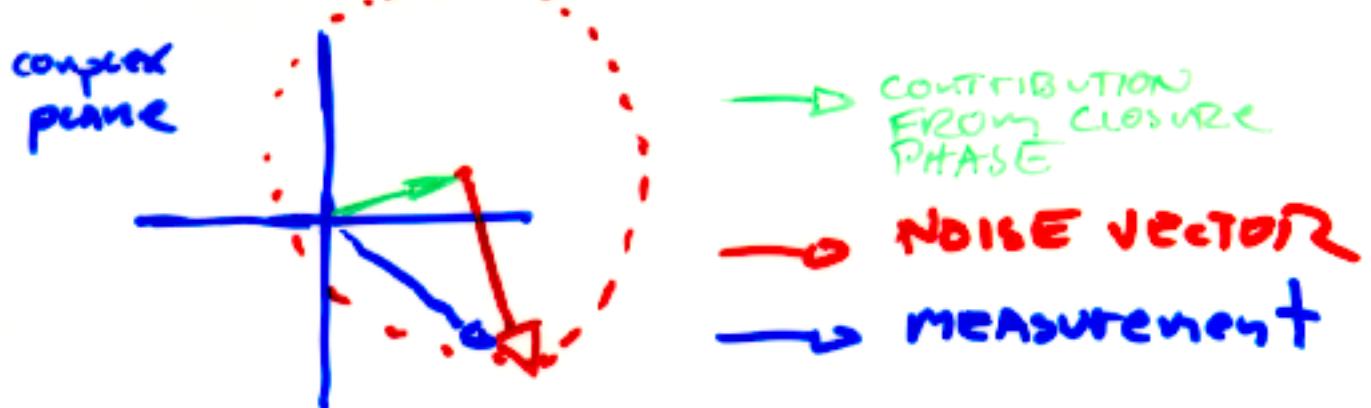
DEFINE A_{ijkl} :

$$A_{ijkl} = \frac{|\gamma_{ik}| |\gamma_{jl}|}{|\gamma_{ik}| |\gamma_{jl}|} = \frac{|G_i| |G_j| |\tilde{\gamma}_{ij}| |G_k| |G_l| |\tilde{\gamma}_{kl}| |\tilde{\gamma}_{jk}^{\text{TRUE}}|}{|G_i| |G_j| |\tilde{\gamma}_{ik}^{\text{TRUE}}| |G_k| |G_l| |\tilde{\gamma}_{jl}^{\text{TRUE}}|}$$

$$= \frac{|\tilde{\gamma}_{ij}^{\text{TRUE}}| |\tilde{\gamma}_{kl}^{\text{TRUE}}|}{|\tilde{\gamma}_{ik}^{\text{TRUE}}| |\tilde{\gamma}_{jl}^{\text{TRUE}}|}$$

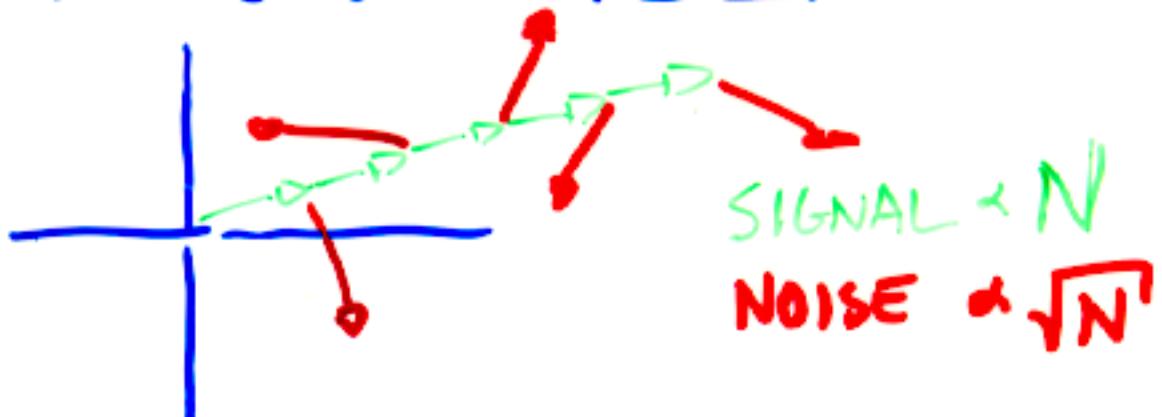
THIS QUANTITY
IS INDEPENDENT
OF TELESCOPE-
SPECIFIC GAINS

CLOSURE PHASE MEASUREMENTS
SUFFER FROM ADDITIVE NOISE,
NOT "PHASE" NOISE —

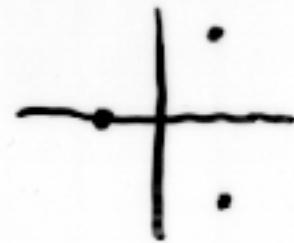
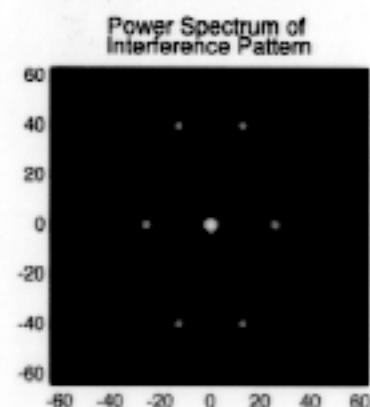
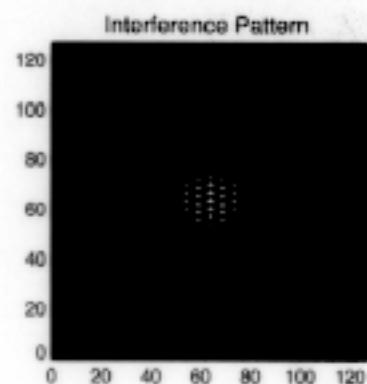
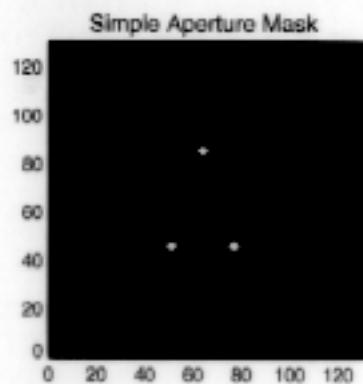
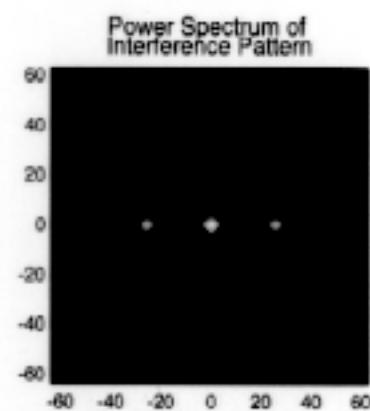
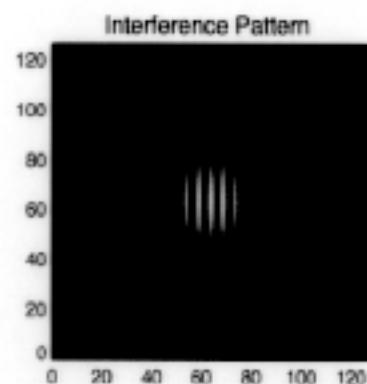
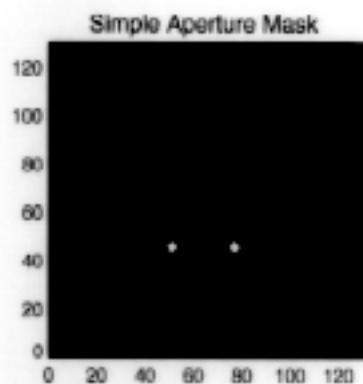


$$\hat{B}_{ijk}^{\text{measured}} = \hat{B}_{ijk}^{\text{TRUE}} + \text{NOISE}$$

BUT, AVERAGING WORKS —



HENCE, SNR OF BISPECTRUM $\propto \sqrt{N}$



THE EFFECT OF THE PHASE CENTER

$$\tilde{\gamma}' \rightarrow \tilde{\gamma} \cdot e^{i \bar{k} \cdot \bar{\Delta x}}$$

Hence,

$$\phi'_{ij} \rightarrow \phi_{ij} + \bar{k}_{ij} \cdot \bar{\Delta x}$$

$$\phi'_{jik} \rightarrow \phi_{jik} + \bar{k}_{jik} \cdot \bar{\Delta x}$$

$$\phi'_{kij} \rightarrow \phi_{kij} + \bar{k}_{kij} \cdot \bar{\Delta x}$$

$$\Phi'_{ijk} \rightarrow \Phi_{ijk} + [\bar{k}_{ij} + \bar{k}_{jik} + \bar{k}_{kij}] \cdot \bar{\Delta x}$$

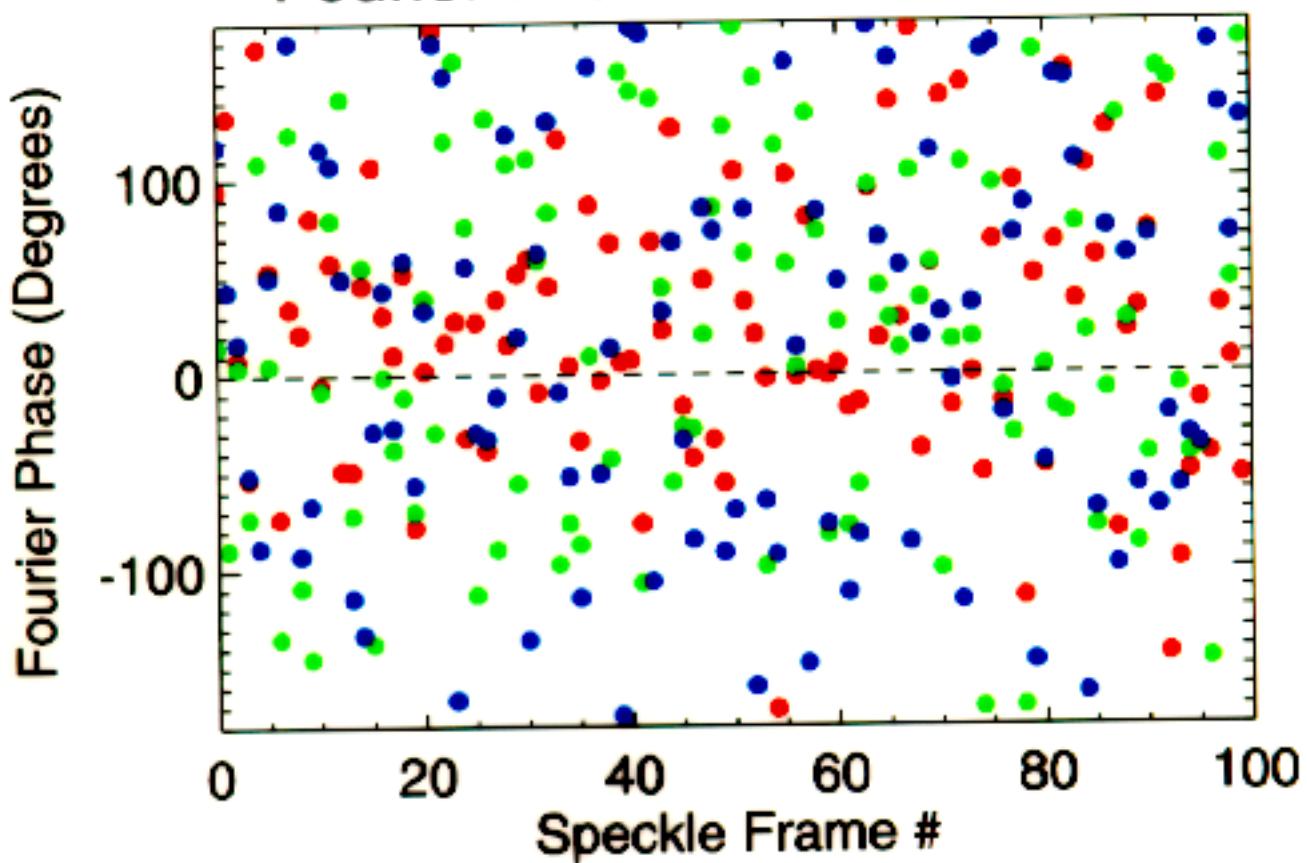
But, what is \bar{k}_{ij} ? $\bar{k}_{ij} = 2\pi \frac{\bar{B}_{ij}}{\lambda}$

Hence, $\bar{k}_{ij} + \bar{k}_{jik} + \bar{k}_{kij} \propto \underbrace{\bar{B}_{ij} + \bar{B}_{jik} + \bar{B}_{kij}}$

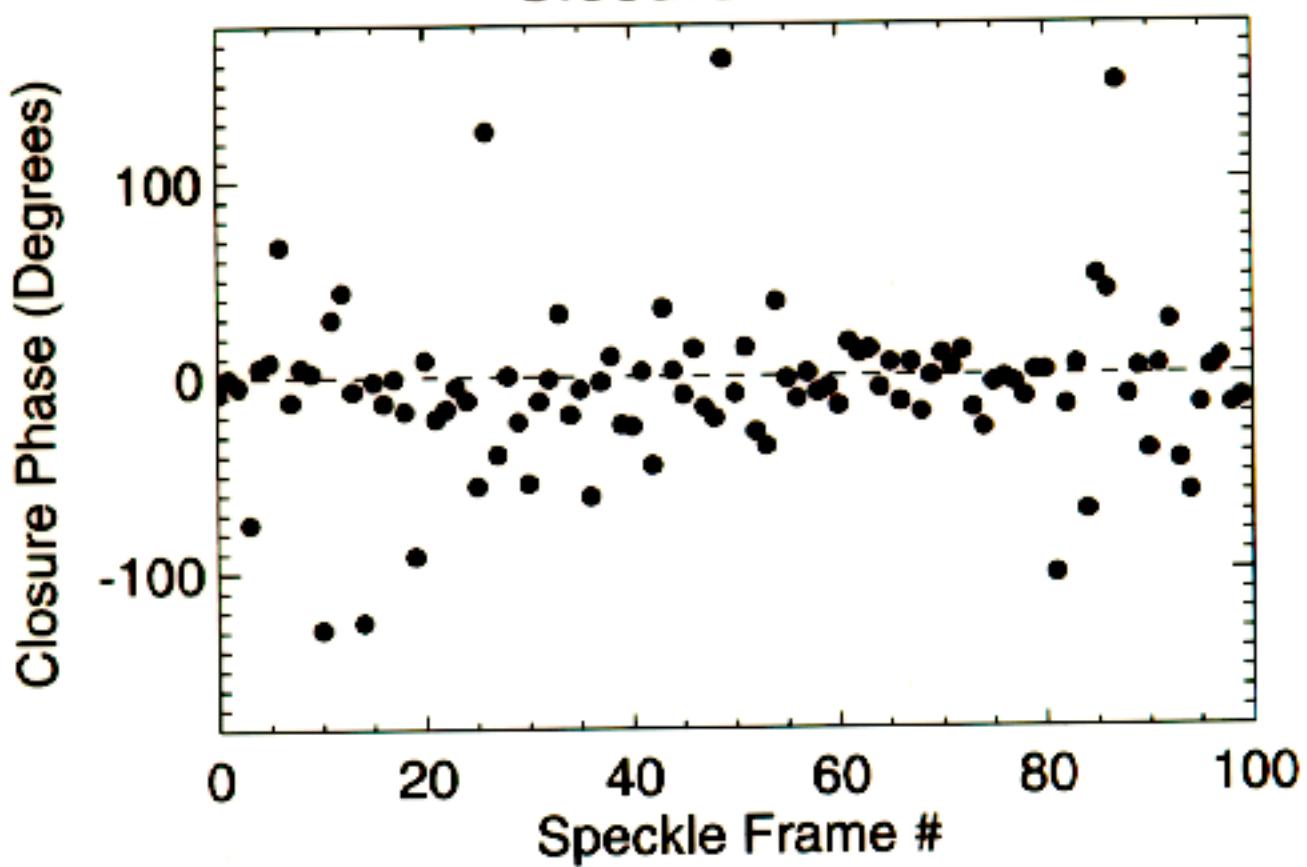


ZERO FOR CLOSING TRIANGLES

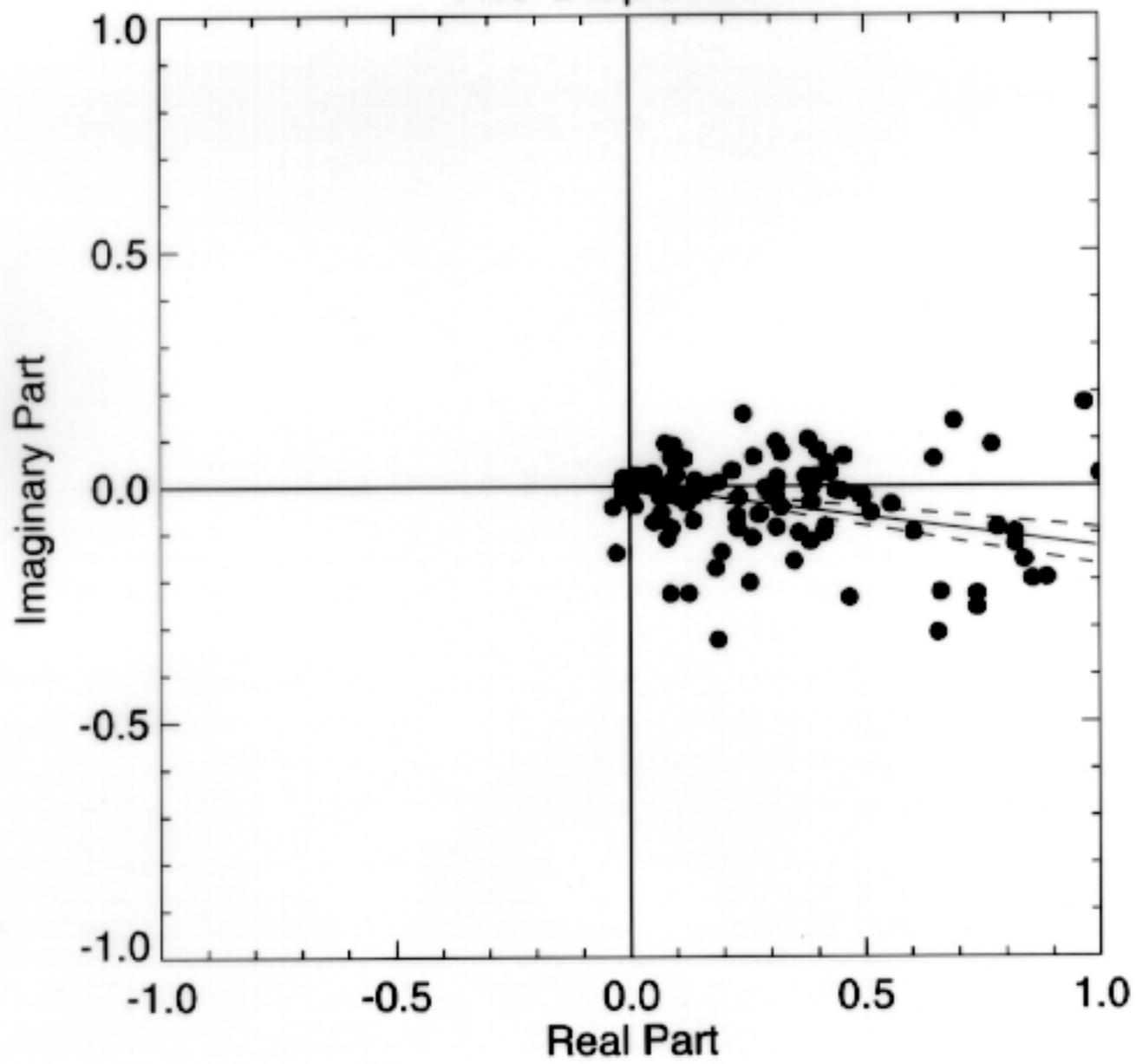
Fourier Phases on 3 Baselines



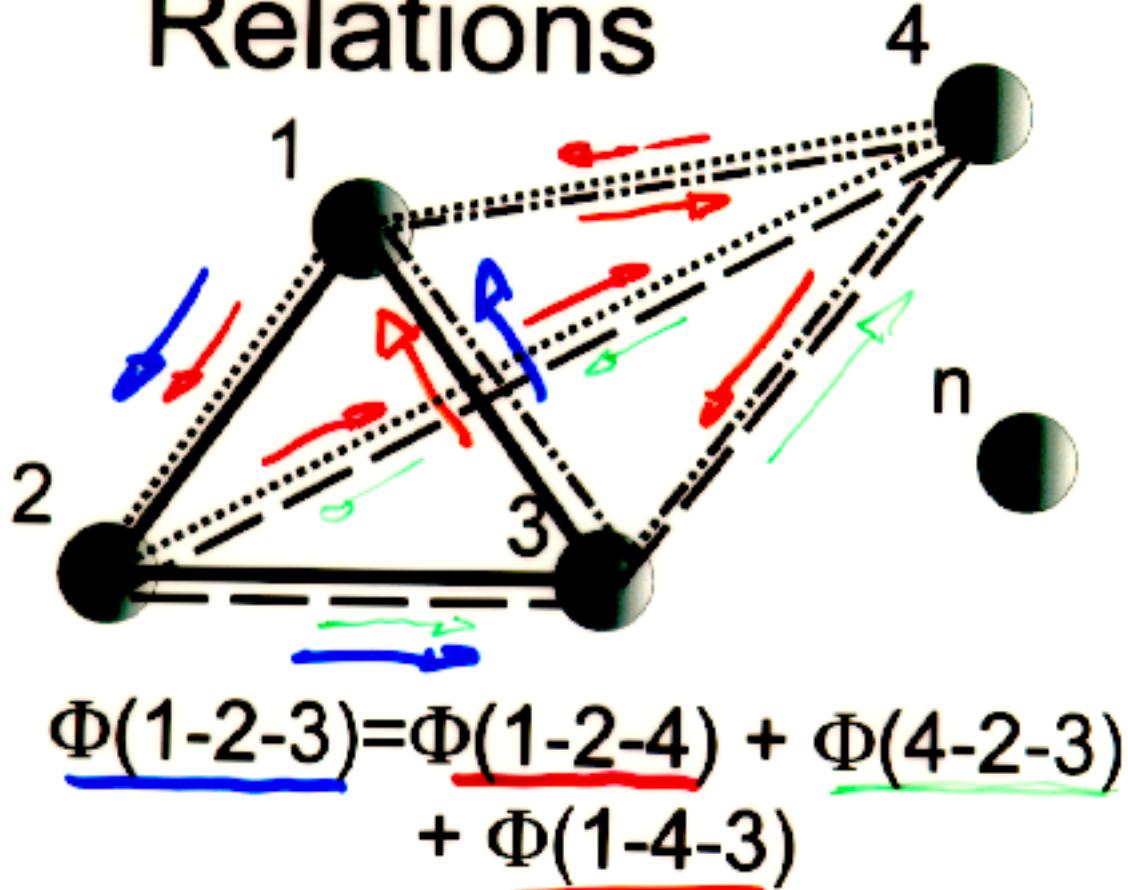
Closure Phase



The Bispectrum



Closure Phase Relations



In General:

$$\Phi(1-2-3) = \Phi(1-2-n) + \Phi(n-2-3) + \Phi(1-n-3)$$

FOR NON-REDUNDANT ARRAY,

NUMBER OF INDEPENDENT $\underline{\Phi}$ CP:

$$\binom{N-1}{2} = \frac{(N-1)(N-2)}{2}$$

NUMBER OF FOURIER PHASES:

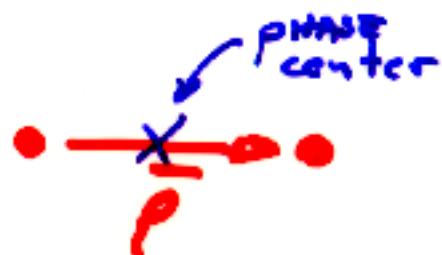
$$\binom{N}{2} = \frac{N \cdot (N-1)}{2}$$

N	NUMBER OF PHASES	NUMBER OF $\underline{\Phi}$ CP	% OF PHASE INFO
3	3	1	33%
7	21	15	71%
21	210	190	90%
27	351	325	93%
50	1225	1176	96%

EXAMPLE: AN EQUAL BINARY

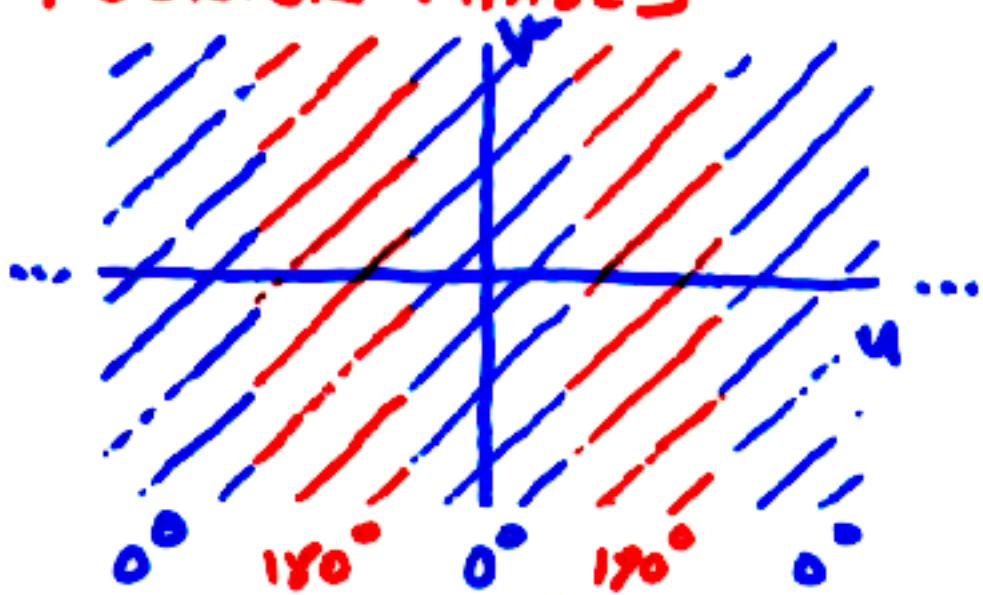
$$\tilde{Y} = e^{-i\frac{\vec{P} \cdot \vec{k}}{2}} + e^{i\frac{\vec{P} \cdot \vec{k}}{2}}$$

$$= 2 \cos\left(\frac{\vec{P} \cdot \vec{k}}{2}\right)$$



Hence, PHASE
EITHER $0^\circ, 180^\circ$

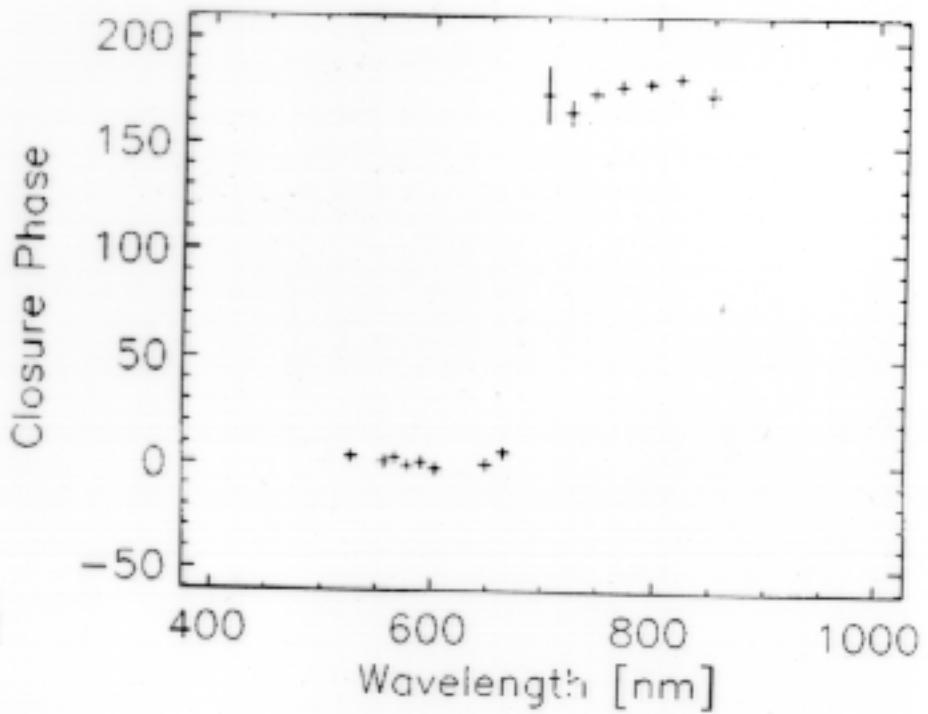
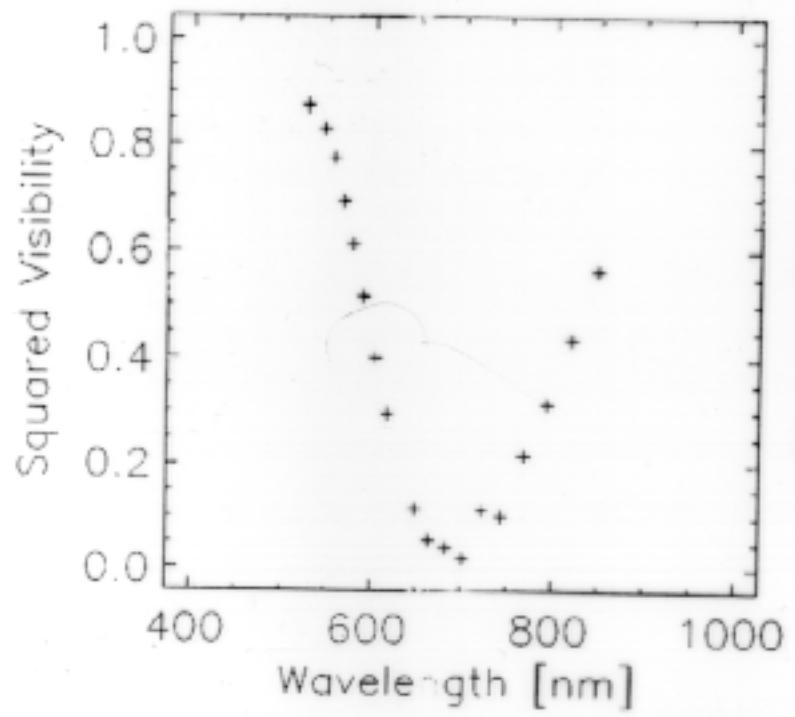
FOURIER PHASES



NOTICE ABRUPT PHASE JUMPS
WHEN VISIBILITY AMPLITUDE
GOES THROUGH NULLS

PHASE DISCONTINUITIES SMOOTHED
OUT FOR UNEQUAL BINARIES

BENSON ET AL.: ξ^1 UMa



EXAMPLE: FAINT HOT SPOT ON STELLAR SURFACE



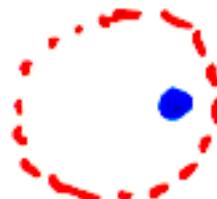
NOTE:
ELLIPTICAL
STARS YIELD
 $\Phi_{CP} = 0^\circ, 180^\circ$
ONLY.

AT LOW SPATIAL RESOLUTION,



STELLAR DISK IS
CENTRALLY SYMMETRIC,
i.e. $\Phi_{CP} = 0^\circ, 180^\circ$

AT HIGH SPATIAL RESOLUTION,



STELLAR DISK IS
RESOLVED, LEAVING
ONLY THE HOT SPOT,
i.e. $\Phi_{CP} = 0^\circ, 180^\circ$

HOWEVER, CLOSING-TRIANGLES OF
INTERMEDIATE BASELINES WILL
YIELD NON-ZERO Closure PHASES

MODELS: DISK + HOTSPOTS

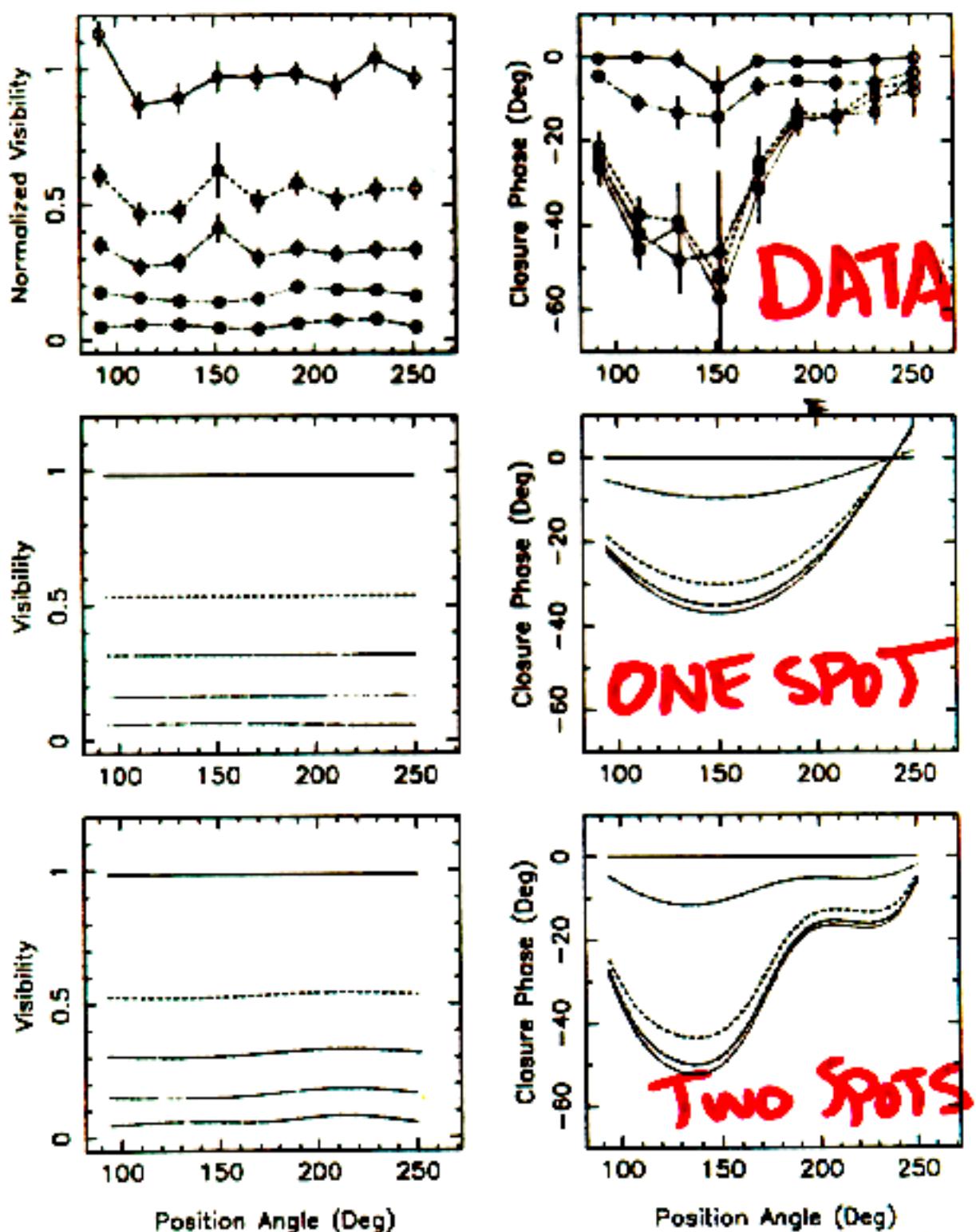
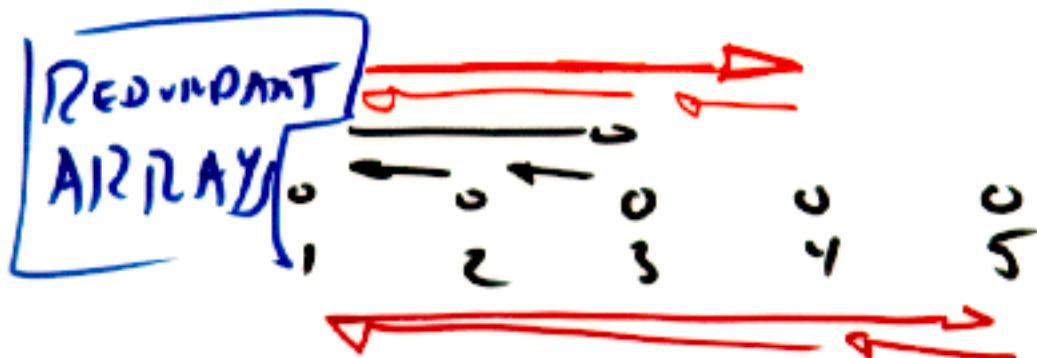


Figure 4. Visibility amplitudes and closure phases as measured

TUTHILL 1994

IMAGING

(CONVERTING RISPECTRUM TO FOURIER PHASES.)



$$\phi_{12} = 0$$

$$q\phi_{13} = \phi_{132} - \cancel{\phi_{32}} - \phi_{11}$$

SOLVE FOR ALL PHASES
DIRECTLY!

BUT! POOR NOISE PROPAGATION
CHARACTERISTICS.

INEFFICIENT USE OF FREQUENCIES
POOR (N, N) COVERAGE.

~~LINEAR~~ LINEAR OPERATION \Rightarrow MATRICES.
EXCEPT COMPLEX LOGIC $\bmod 2\pi$.

Image Reconstruction

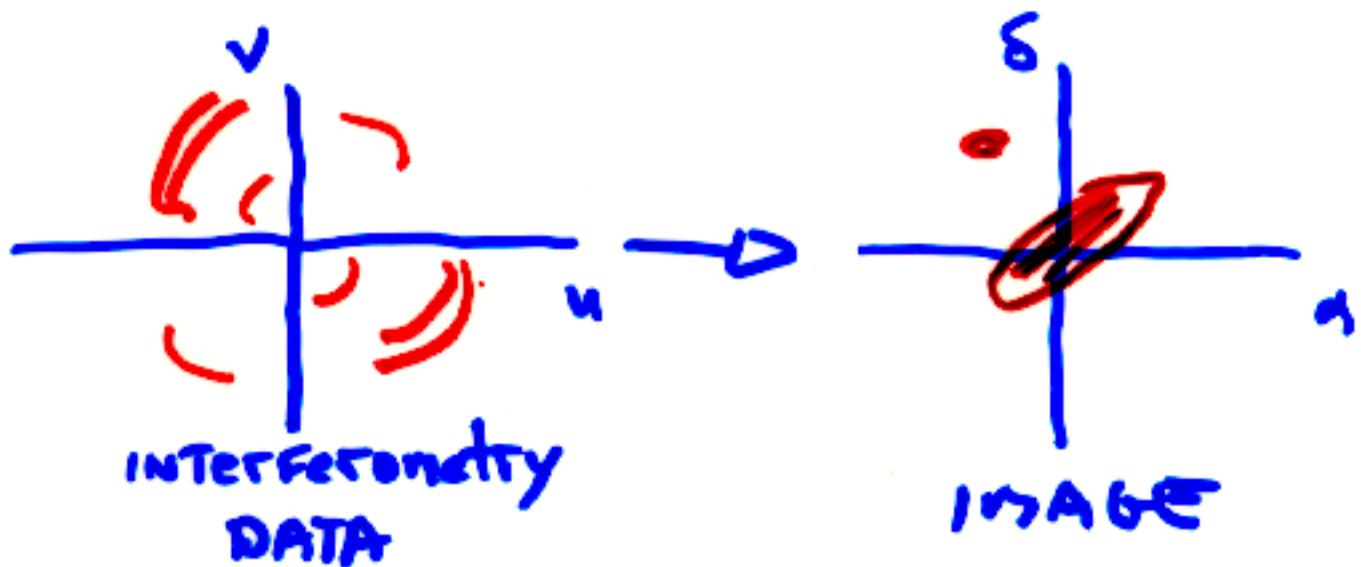
- Find image which fits both visibility amplitudes and closure phases (to within experimental uncertainty)

Further Constraints

- Positive Definite
- Limited Field-of-View
- "Smooth" – e.g. MEM
- Use *a priori* information

"STANDARD" APERTURE SYNTHESIS IMAGING:

converts unevenly SAMPLED
FOURIER PHASES AND AMPS
(i.e. complex VISIBILITY) INTO
A MAP FREE OF ARTIFACTS



ALGORITHMS: CLEAN
MEM

Iteration Cycle

Start with Initial Model



**Generate Fourier phases
consistent with Closure
phases, using Model phases
as necessary.**



**Using these candidate phases
and visibility amplitudes, use
CLEAN/MEM to generate map**



**If not converged, use new map
as starting Model and repeat**

Self-Calibration

Models intrinsic Fourier
phases *plus* telescope errors

Generate Fourier phases
consistent with Closure Phases &
begin with initial trial Image



Calculate complex Visibility
(Amplitudes & Phases) of trial image



Adjust telescope errors so "measured"
Fourier phases are best fit by
combination of trial image phases plus
telescope errors



Correct trial phases based on new
estimates of telescope errors, and
map using CLEAN/MEM



If not converged, use new map
as the next trial image

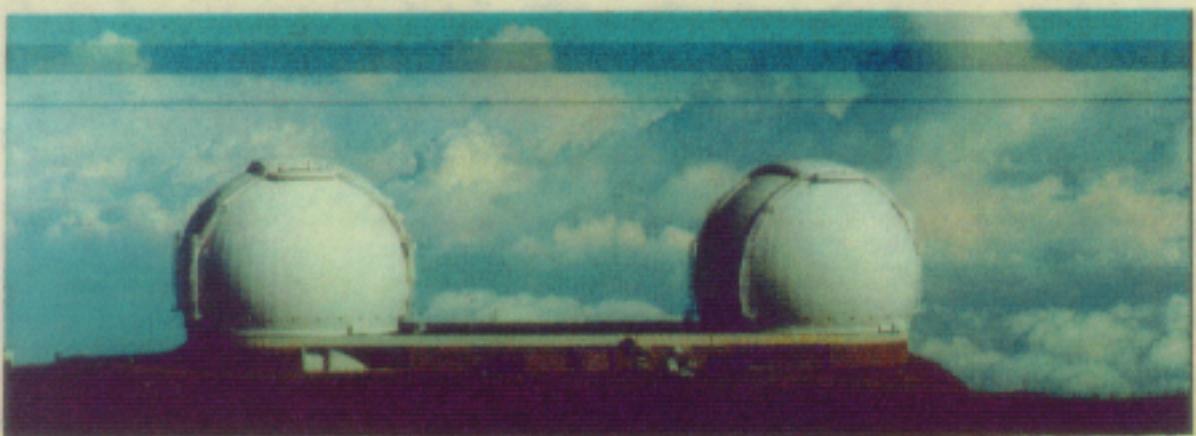
$$\underline{\Phi}_{\text{"MEASURED"}} = \underline{\Phi}_{\text{TRUE}} + [\phi_j - \phi_i]$$

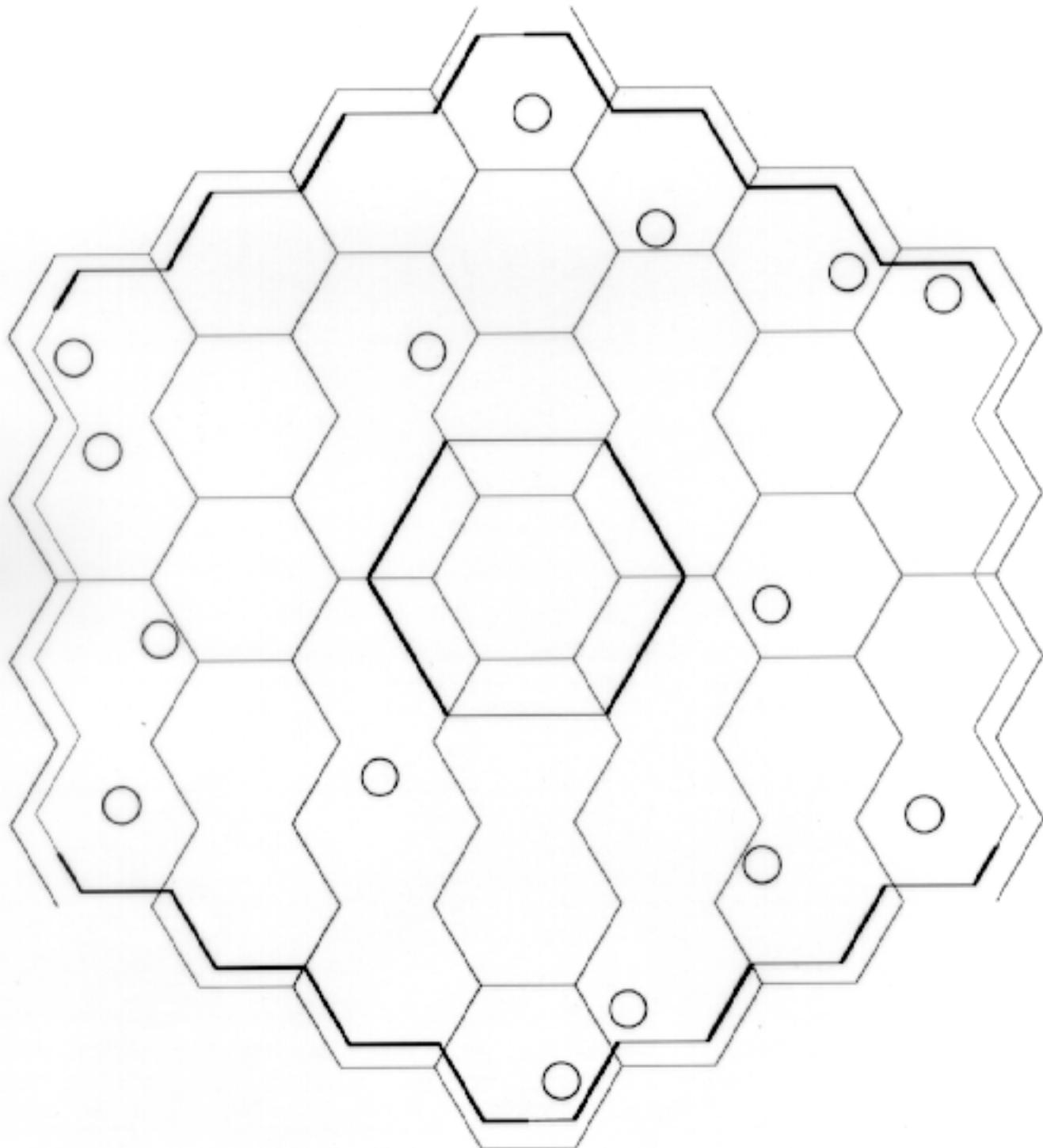

ATMOSPHERE
 DOES NOT
 CHANGE THE
 CLOSURE PHASES

ADJUST THESE
 TERMS ITERATIVELY,
 APPLYING IMAGE
 CONSTRAINTS DURING
 EACH CYCLE.

Works extremely well –
 FOR HIGH SNR (~ 5)
 AND GOOD FOURIER
 COVERAGE.

EMERGING IMAGING INTERFEROMETERS
 (3-6 telescopes) PROBING
 NEW REGIME...





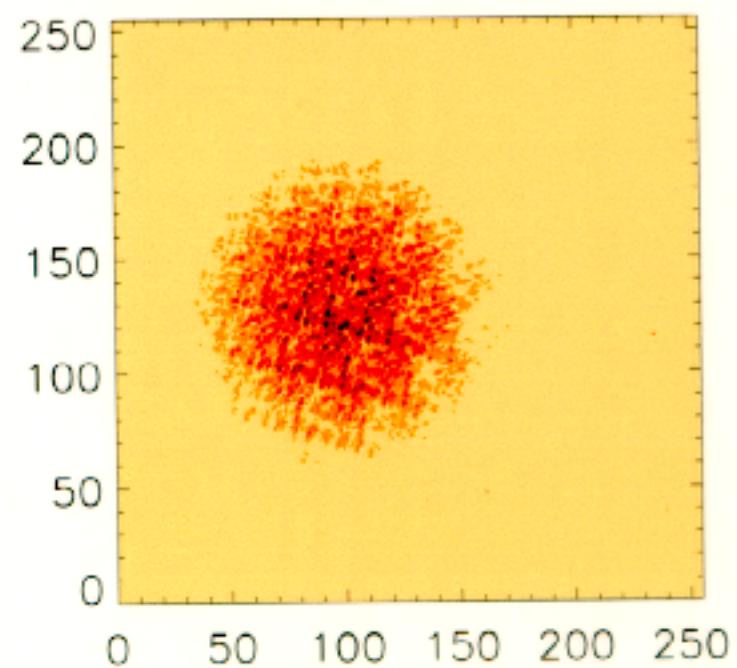
Keck Pupil Mask

DRAWN BY: Peter Tuthill DATE: 25-Feb-97

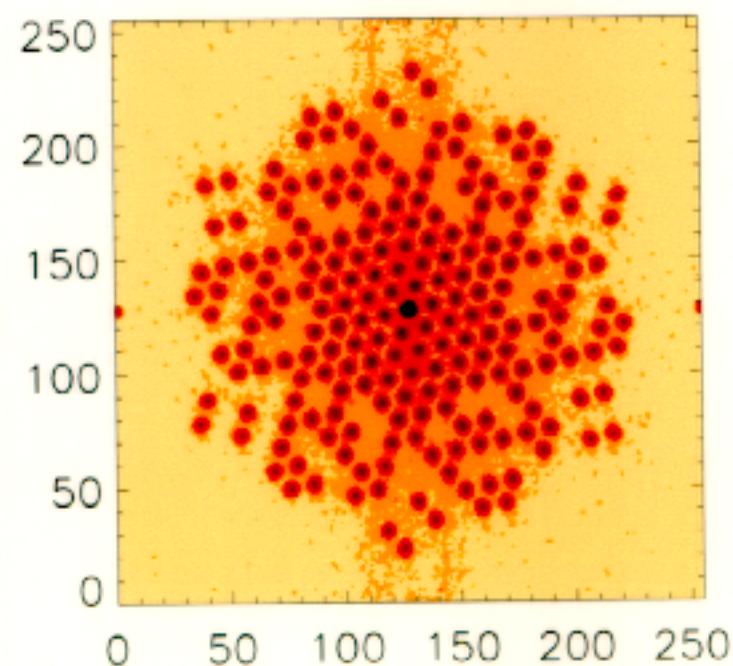
SCALE: 1:50 MATERIAL: Aluminum

15 Hole Golay (35cm holes)

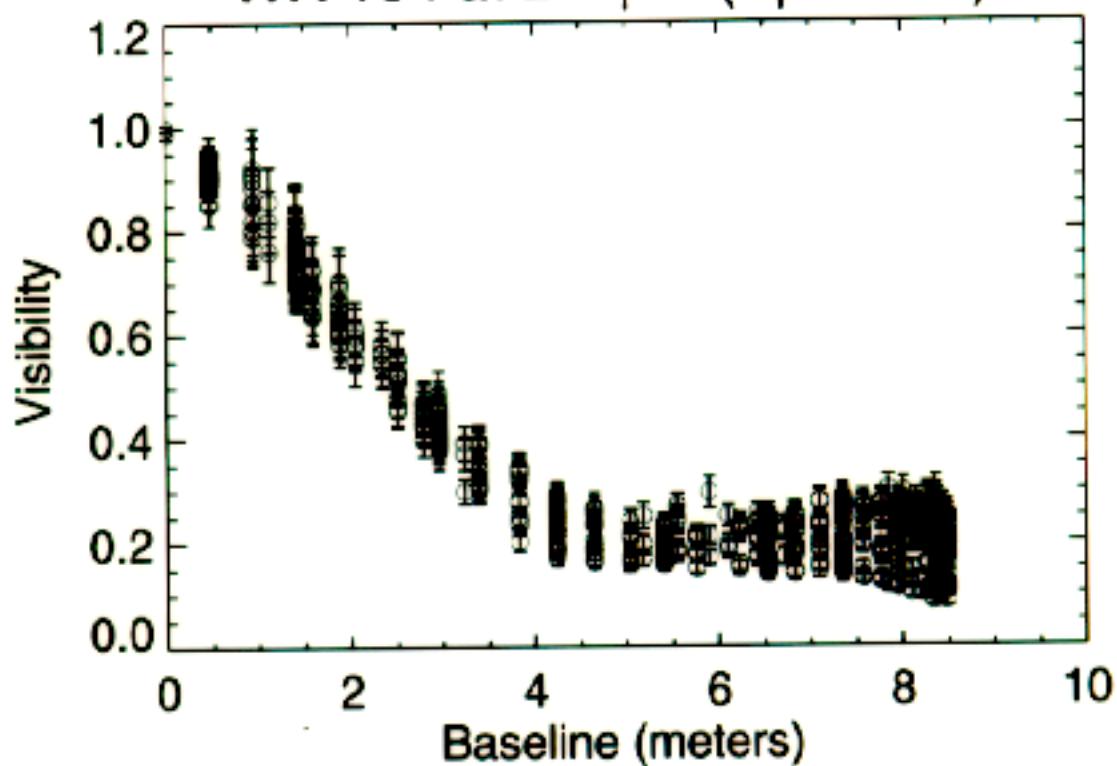
Speckle Pattern



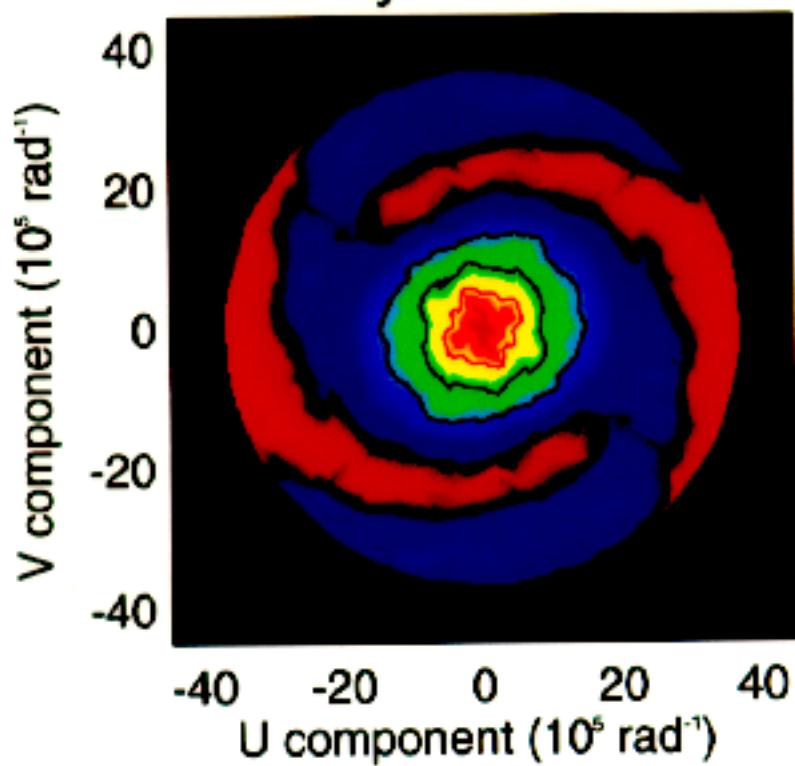
Power Spectrum



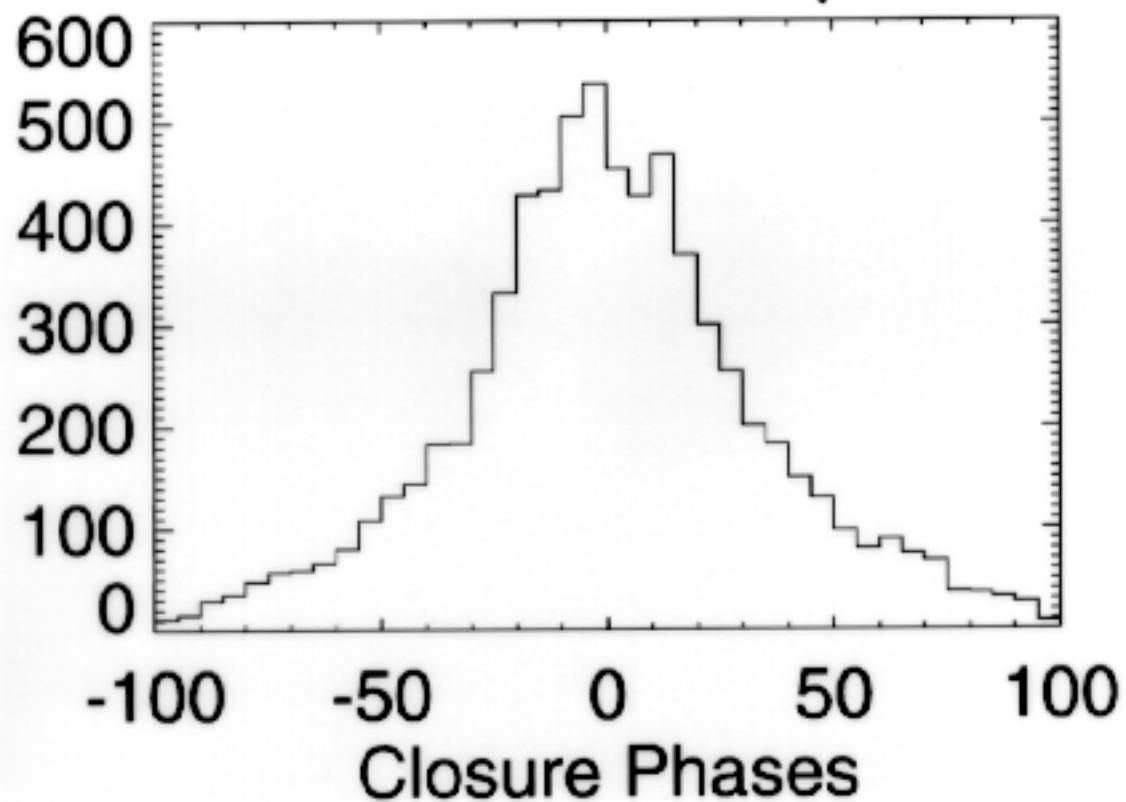
WR 104 at 2.2 μ m (Apr 1998)



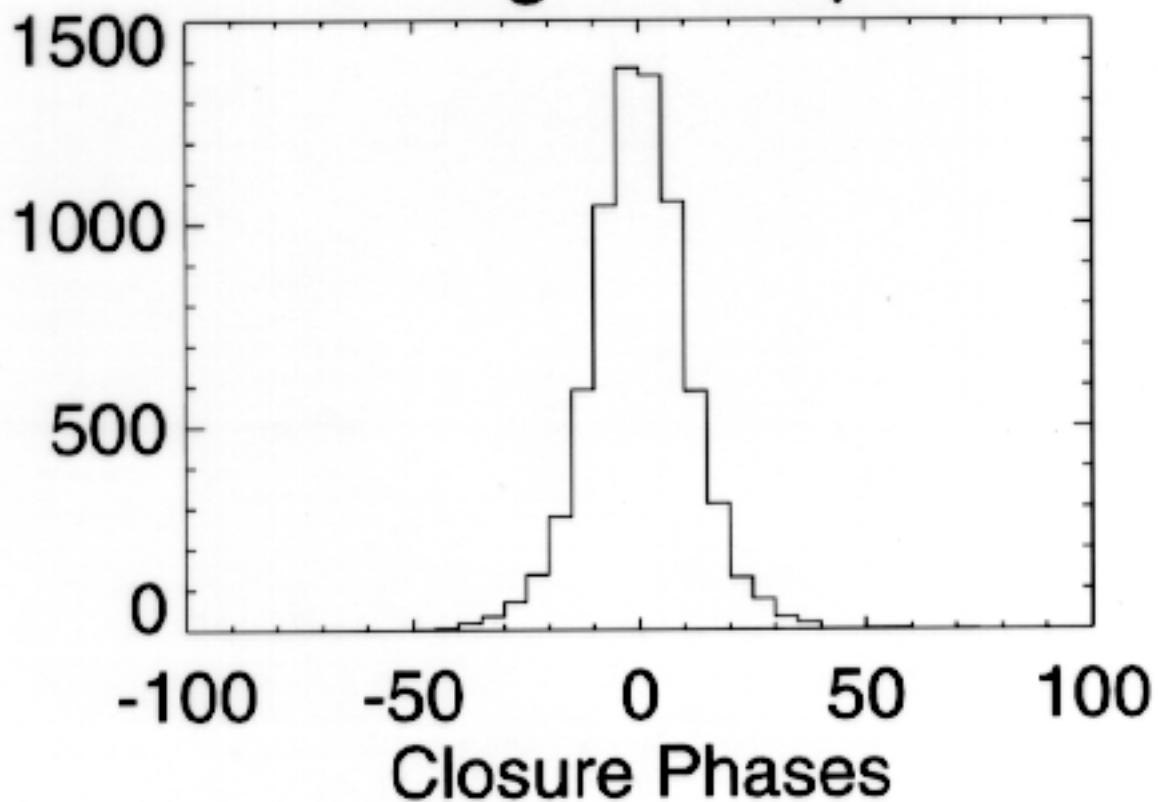
Visibility in U-V Plane



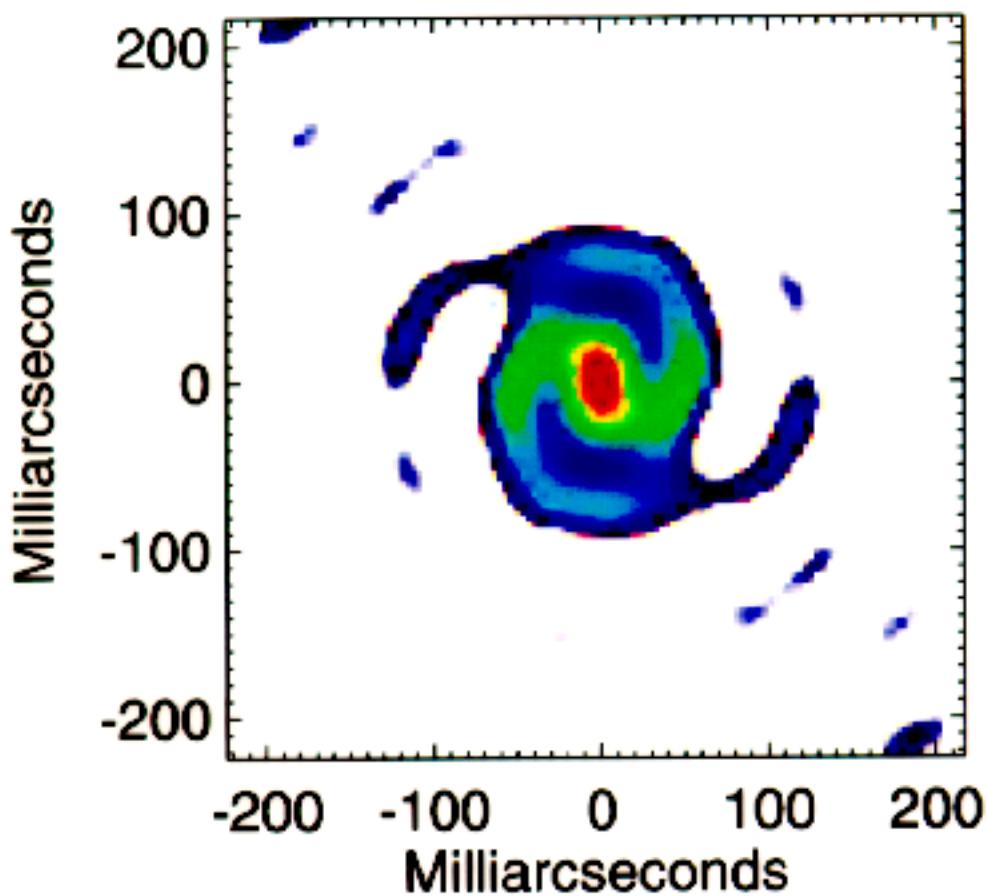
WR104 at 2.2 μ m



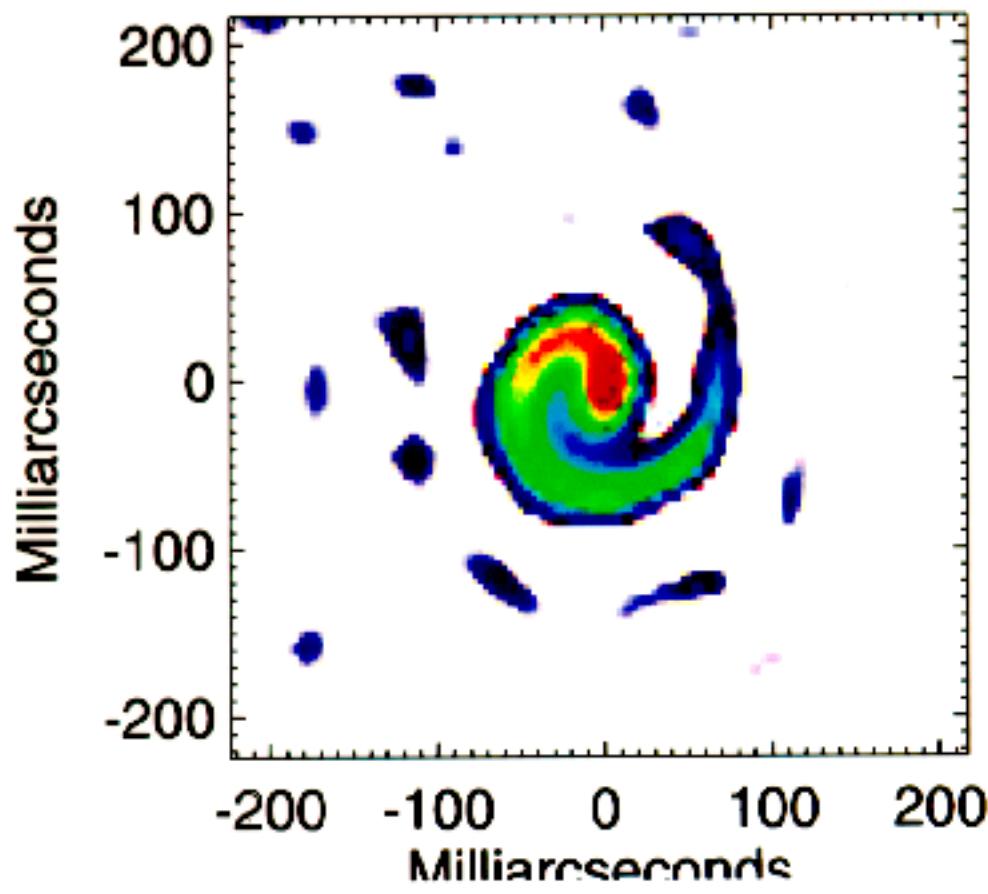
14 Sgr at 2.2 μ m



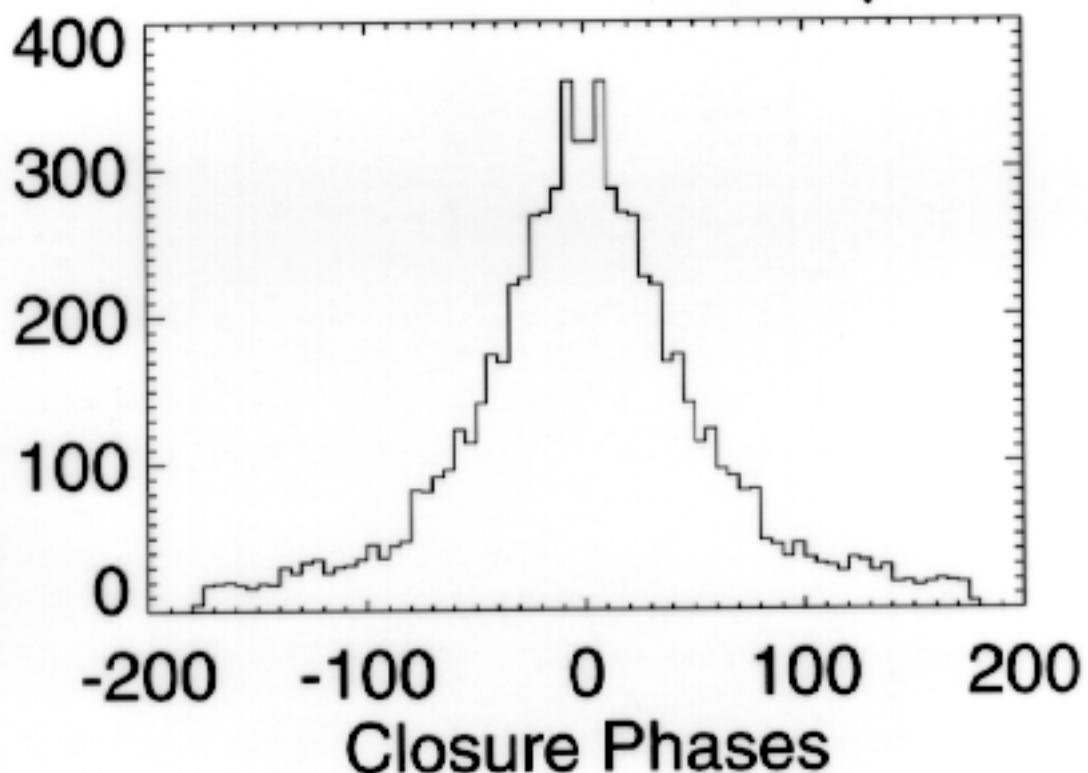
WR 104 w/o Closure Phases



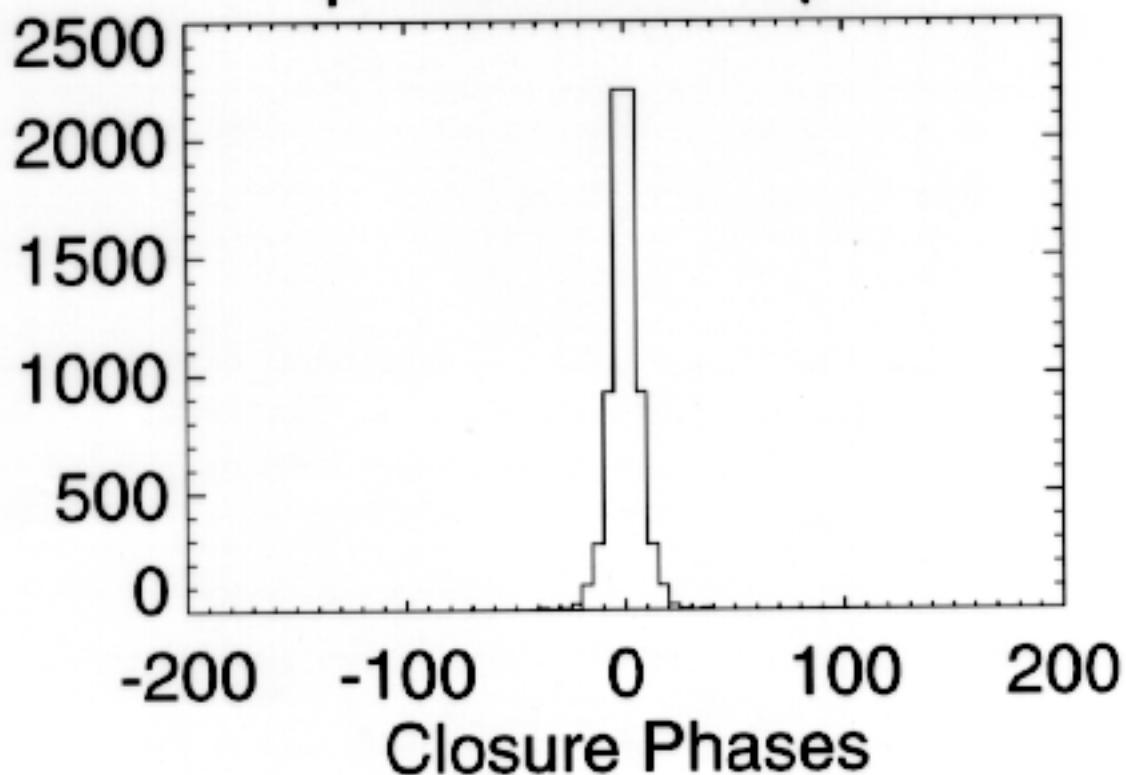
WR 104 with Closure Phases



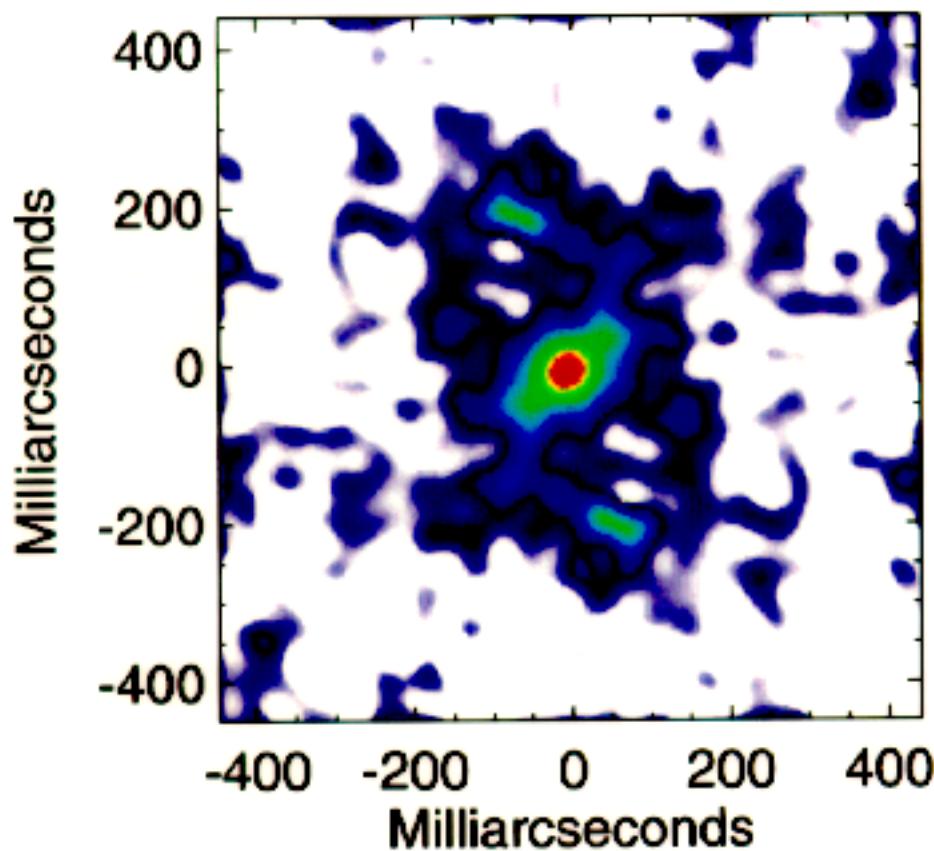
IRC+10216 at $2.2\mu\text{m}$



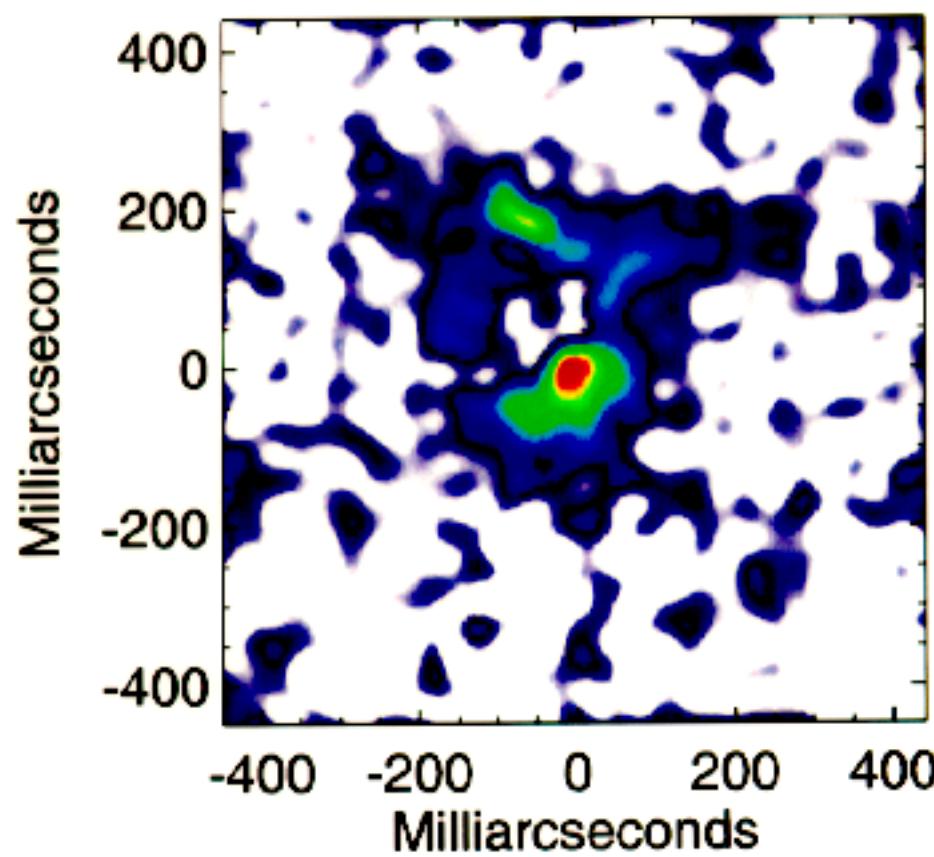
π Leo at $2.2\mu\text{m}$



IRC+10216 w/o Closure Phases



IRC+10216 with Closure Phases



The Future

Outstanding Problems with Mapping

- Should fit directly to closure phases;
outside of self-cal paradigm
(computationally expensive but
straightforward)
- Baseline-dependent residuals (limitation
of χ^2 ; use multi-resolution approach)
- Point source(s) embedded in extended
nebulosity (model/mapping hybrids,
i.e. use *a priori* info)
- Artifacts of (highly) uneven Fourier
coverage

Which Triangles?

- Given a sparse array, which closure
phases are most important to measure?
- Better understanding of how closure
phases constrain the parameter space of
possibility Fourier phases

Important Points

- Closure phases (and amplitudes) are insensitive to telescope-specific errors such as atmospheric fluctuations
- They can not be used to calibrate baseline-dependent problems
- Non-zero closure phases indicate asymmetric structure; however, point-symmetric objects (such as disks) have Φ_{CP} of 0° or 180°
- Software and theoretical work necessary for optimal imaging with the latest generation of interferometers
- Despite the total scrambling of Fourier phases by a turbulent atmosphere, the use of closure phases and image constraints allow nearly complete recovery of all phase information